# Birth-and-death evolutions in random environments

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### What is modelled?

Tumour evolution and related quantities such as

- Volume/size of tumour
- Growth of tumour patters, speed of growth, shape of growth
- Mutations, heterogenity and homogenity
- Interplay between the tumour and
  - Immune system
  - Healthy cells
  - Environmental factors
- And many other ...

Models should be as simple as possible, but as detailed as necessary. Simplification: Tumour can be described on different scales

### microscopic < mesoscopic < macroscopic

### How do we model?

Ordinary differential equations (also systems of ODEs)

- Mesoscopic or macroscopic desription.
- Describe specific quantities of the tumour in certain regimes.
- $\bullet\,$  Comparably easy for the analysis and simulations  $\rightarrow\,$  compare with data.
- In many cases deterministic.
- No spatial structure.
- How to justify such equations?

What are the right equations?

- Stochastic models
  - microscopic description of cells, mutations, interactions
  - Describe collection of cells as a stochastic process.
  - Analysis and simulations are more challenging
    - $\rightarrow$  how to compare with experimental data?
  - Include a spatial structure.

What are the right models?

### What are we interested in?

- O Description of microscopic tumours by spatial birth-and-death models.
- 2 Reducing complexity
  - What are the main building blocks?
  - Derivation of effective equations from these models  $\rightarrow$  spatial analogues of ODEs
  - What is the effect of: immune system, other cells, environmental factors?
- **③** Time evolution of (spatial) correlations
- Invariant states, equilibrium states
- What is not done: Comparison with data

### What could/should, in principle, be done

- Cells with additional marks (mutations, fitness, ...)
- Ø Motion of cells (migration, metastasis, go-or-growth, ...)
- **(a)** Non-Markov dynamics (equations with delay, fractional derivatives in time, ...)
- Time-inhomogeneous models (time-dependent parameters due to therapy)

# System of Tumour cells

We suppose that

- Tumour cells are indistinguishable (among one type)
- Sufficiently many cells, i.e. statistical description is adequate.
- Each cell has position and maybe other traits.
   It can be represented as an element x ∈ ℝ<sup>d</sup>.

The collection of all cells forms a microscopic state

$$\gamma = \{ x_n \in \mathbb{R}^d \mid n \ge 1 \}.$$

System may be:

- finite, i.e.  $|\gamma| < \infty$ .
- infinite, i.e.  $|\gamma \cap K| < \infty$  for all balls  $K \subset \mathbb{R}^d$ .

$$\Gamma^{\mathcal{S}} = \{ \gamma \subset \mathbb{R}^d \mid |\gamma \cap \mathcal{K}| < \infty \text{ for all balls } \mathcal{K} \subset \mathbb{R}^d \}.$$

We mainly focus on the infinite case.

Environment is another particle system with microscopic state

$$\omega = \{ y_n \in \mathbb{R}^d \mid n \ge 1 \}.$$

Configuration space

$$\Gamma^{E} = \{ \omega \subset \mathbb{R}^{d} \mid |\omega \cap K| < \infty \text{ for all balls } K \subset \mathbb{R}^{d} \}.$$

Environment could be

- Fixed configuration  $\omega \in \Gamma^{E}$ .
- An equilibrium process on  $\Gamma^E$  with some invariant measure.
- Non-equilibrium spatial birth-and-death process.

Joint microscopic state is  $(\gamma, \omega)$ , i.e. an element in

 $\Gamma^2 := \Gamma^S \times \Gamma^E.$ 

# Statistical description

Observable quantities are

$$\langle F \rangle_{\mu} := \int_{\Gamma^{S} \times \Gamma^{E}} F(\gamma, \omega) \mathrm{d}\mu(\gamma, \omega)$$

where  $\mu$  is a probability measure (= state) on  $\Gamma^{\rm S} \times \Gamma^{\rm E}$ 

• Number of tumour cells in volume  $\Lambda^{S}$ 

$$\int_{\Gamma^{\mathcal{S}}} |\gamma \cap \Lambda^{\mathcal{S}}| \mathrm{d}\mu(\gamma, \omega).$$

 $\bullet$  Second order correlations:  $\Lambda^{\mathcal{S}} \times \Lambda^{\mathcal{E}} \subset \mathbb{R}^d \times \mathbb{R}^d$ 

$$\int_{\Gamma^S \times \Gamma^E} |\gamma \cap \Lambda^S| |\omega \cap \Lambda^E| \mathrm{d}\mu(\gamma, \omega).$$

• Higher order correlations

$$\int_{\Gamma^{S}\times\Gamma^{E}}\prod_{k=1}^{n}|\gamma\cap\Lambda_{k}^{S}|\cdot\prod_{k=1}^{m}|\gamma\cap\Lambda_{k}^{E}|\mathrm{d}\mu(\gamma,\omega).$$

# Space of admissible measures

 $\mu$  state on  $\Gamma^2$ . Correlation function  $k_{\mu}^{(n,m)}$  is locally integrable function s.t.

$$\int_{\Gamma^{S}\times\Gamma^{E}}\prod_{k=1}^{n}|\gamma\cap\Lambda_{k}^{S}|\cdot\prod_{k=1}^{m}|\gamma\cap\Lambda_{k}^{E}|\mathrm{d}\mu(\gamma,\omega)=\frac{1}{n!}\frac{1}{m!}\int_{\Lambda_{1}^{+}\times\cdots\times\Lambda_{n}^{+}}\int_{\Lambda_{1}^{-}\times\ldots\Lambda_{m}^{-}}k_{\mu}^{(n,m)}\mathrm{d}x\mathrm{d}y.$$

Example:

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• 
$$k^{(0,0)} = \mu(\Gamma^2) = 1.$$
  
•  $n + m = 1$  yields for  $\Lambda \subset \mathbb{R}^d$  compact  

$$\int_{\Gamma^2} |\gamma^+ \cap \Lambda| \mathrm{d}\mu(\gamma) = \int_{\Lambda} k^{(1,0)}_{\mu}(x) \mathrm{d}x, \quad \int_{\Gamma^2} |\gamma^- \cap \Lambda| \mathrm{d}\mu(\gamma) = \int_{\Lambda} k^{(0,1)}_{\mu}(y) \mathrm{d}y.$$

Up to some mathematical aspects, we obtain a one-to-one correspondence

states 
$$\mu \longleftrightarrow$$
 correlation functions  $k_{\mu} = (k_{\mu}^{(n,m)})_{n,m=0}^{\infty}$ 

## Dynamics

Model time evolution by an evolution of states

$$\mu_0\longmapsto \mu_t \quad \text{ or } \quad k^{(n,m)}_{\mu_0}\longmapsto k^{(n,m)}_{\mu_t}.$$

Elementary events:

• Birth: Each particle creates a new particle, a particle may appear from the outside.

$$\gamma\longmapsto \gamma\cup\{x\}, \ x\not\in\gamma.$$

Birth rate:  $b(x, \gamma, \omega) \ge 0$ .

• Death: Particles have a lifetime and compete for resources.

$$\gamma \longmapsto \gamma \setminus \{x\}, \ x \in \gamma.$$

Death rate:  $d(x, \gamma, \omega) \ge 0$ .

• Motion, migration, mutations and many others are possible.

 $\omega \in \Gamma^{E}$  takes influence of environment into account.

Let  $L^{S}_{\gamma}(\omega)$  be the Markov operator for the tumour cells, i.e.

$$(L^{S}_{\gamma}(\omega)F)(\gamma) = \sum_{x \in \gamma} d(x, \gamma \setminus x, \omega)(F(\gamma \setminus x) - F(\gamma)) + \int_{\mathbb{R}^{d}} b(x, \gamma, \omega)(F(\gamma \cup x) - F(\gamma)) dx.$$

Environment is assumed to be of similar form but with different birth-and-death rates. These rates will be specified later on.

## Dynamics

Markov evolution is described by solutions to (backward) Kolmogorov equation

$$\frac{\partial F_t}{\partial t} = (L_{\gamma}^{S}(\omega) + L_{\omega}^{E})F_t, \quad F_t|_{t=0} = F.$$

Then  $\langle F_t \rangle_{\mu}$  is the time evolution of the expected value of F in state  $\mu$ .

We are interested in the evolution of states  $\mu_0 \mapsto \mu_t$ . We should have  $\langle F_t \rangle_{\mu} = \langle F \rangle_{\mu_t}$ .

Can rewrite this into a system of equations

$$\frac{\partial k_t^{(n,m)}}{\partial t} = (L^{\Delta} k_t)^{(n,m)}, \ k_t^{(n,m)}|_{t=0} = k_0^{(n,m)}$$

where  $L^{\Delta}$  is a double-matrix. Then

$$(k_t^{(n,m)})_{n,m=0}^{\infty} \longleftrightarrow \mu_t.$$

# Evolution of correlation functions

$$\frac{\partial k_t^{(n,m)}}{\partial t} = (L^{\Delta} k_t)^{(n,m)}, \quad k_t^{(n,m)}|_{t=0} = k_0^{(n,m)}$$

- $L^{\Delta}$  can, for many models, be computed explicitly from  $L^{S}_{\gamma}(\omega), L^{E}_{\omega}$ .
- Explicit form in coordinates

$$\frac{\partial k_t^{(n,m)}}{\partial t} = \sum_{k,l=0}^{\infty} L_{kl,nm}^{\Delta} k_t^{(k,l)}, \quad k_t^{(n,m)}|_{t=0} = k_0^{(n,m)}.$$

• In some cases the right-hand side has recursive structure, i.e.

$$\frac{\partial k_t^{(n,m)}}{\partial t} = \sum_{k,l=0}^{n,m} L_{kl,nm}^{\Delta} k_t^{(k,l)}, \quad k_t^{(n,m)}|_{t=0} = k_0^{(n,m)}.$$

Hence it may be solved explicitly.

**()** We are interested in the projections  $\mu_t^S$  of  $\mu_t$  onto  $\Gamma$ , i.e.

$$\mu_t^{\mathcal{S}}(\mathcal{A}) := \mu_t(\mathcal{A} imes \Gamma^{\mathcal{E}}) \longleftrightarrow k_{\mu_t^{\mathcal{S}}}^{(n)} = k_{\mu_t}^{(n,0)}.$$

Equation for  $k_{\mu_t}^{(n)}$  depends on all  $k_{\mu_t}^{(n,m)}$ . Projection is not Markov.

- Find a closed equation (after proper scaling) for  $\mu_t^S$  or  $k_{\mu_s^S}$ .
- The limiting equation should recover the Markov property.
- Works in a certain regime of parameters on the interactions.

**2** Find closed equations for particle densities  $k_t^{(1,0)}, k_t^{(0,1)}$ , i.e.

$$\frac{\partial \rho_t^{\mathsf{S}}(x)}{\partial t} = v^{\mathsf{S}}(\rho_t^{\mathsf{S}}, \rho_t^{\mathsf{E}})(x), \quad \frac{\partial \rho_t^{\mathsf{E}}(x)}{\partial t} = v^{\mathsf{E}}(\rho_t^{\mathsf{E}})(x).$$

Mesoscopic equations which are obtained after certain scalings.

## Environment

Environment is Glauber dynamics with formal Markov operator

$$(L_{\omega}^{E}F)(\omega) = \sum_{x \in \omega} (F(\omega \setminus x) - F(\omega)) + z \int_{\mathbb{R}^{d}} e^{-E_{\varphi}(x,\omega)} (F(\omega \cup x) - F(\omega)) dx$$

- $z \ge 0$  activity parameter.
- Relative energy

$$E_{arphi}(x,\omega)=\sum_{y\in\omega}arphi(x-y), \;\; x\in \mathbb{R}^d, \;\; \omega\in \Gamma^E.$$

- Death rate is constant to 1.
- Birth rate is given by  $ze^{-E_{\varphi}(x,\omega)}$ .

## Environment

Assumptions

• Interaction potential  $arphi(x)=arphi(-x)\geq 0$  with integrability condition

$$eta(arphi):=\int\limits_{\mathbb{R}}(1-e^{-arphi(x)})\mathrm{d}x<\infty.$$

• Small activity regime

$$z < rac{1}{eeta(arphi)}.$$

Then

- There exists an evolution of states  $(\mu_t)_{t\geq 0}$ .
- There exists a unique invariant measure (Gibbs measure)  $\mu_{\mathrm{inv}}$ .
- Evolution of states is ergodic, i.e.

$$\mu_t \longrightarrow \mu_{\mathrm{inv}}$$
 or  $k_{\mu_t}^{(n)} \longrightarrow k_{\mu_{\mathrm{inv}}}^{(n)}, \forall n, t \longrightarrow \infty.$ 

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Free branching with rates

$$egin{aligned} d(x,\gammaackslash x,\omega) &= m + g\sum_{y\in\omega} d(x-y) \ b(x,\gamma,\omega) &= \sum_{y\in\gamma} a^+(x-y) \end{aligned}$$

- $m \ge 0$  mortality rate of cells.
- $a^+(x-y) = a^+(y-x) \ge 0$  integrable and bounded, proliferation kernel for cells
- $d(x y) = d(y x) \ge 0$  integrable and bounded, interaction with environment.
- $g \ge 0$  coupling constant for interaction with environment.

Consider finite system such that  $m < \lambda := \int a^+(x) dx$ .

#### Free branching process

## Reduced description

Scaling Markov operator  $L^{S}_{\gamma}(\omega) + \frac{1}{\varepsilon}L^{E}_{\omega}$  for  $\varepsilon > 0$  yields when  $\varepsilon \to 0$  reduced description:

$$\overline{d}(x,\gamma\backslash x) = m + \overline{g}(x)$$
  
 $\overline{b}(x,\gamma) = \sum_{y\in\gamma} a^+(x-y)$ 

where

$$\begin{split} \overline{g}(x) &= g \int_{\Gamma^E} \sum_{y \in \omega} d(x - y) \mathrm{d}\mu_{\mathrm{inv}}(\omega) = gz \int_{\Gamma^E} \int_{\mathbb{R}^d} d(x - y) e^{-E_{\varphi}(y,\omega)} \mathrm{d}y \mathrm{d}\mu_{\mathrm{inv}}(\omega) \\ &= g \int_{\mathbb{R}^d} d(x - y) k_{\mu_{\mathrm{inv}}}^{(1)}(y) \mathrm{d}y. \end{split}$$

Consequences

- **()** System is effectively a free branching process with modified mortality rate.
- Space inhomogeneous death rate may be a consequence of interactions with environment.
- O Different environments yield the same reduced description:
  - depends only on invariant state and interactions.
- Invironment may regulate the system:
  - Without environment or interactions the number of particles grows exponentially, since

$$m < \lambda := \int\limits_{\mathbb{R}^d} a^+(x) \mathrm{d}x.$$

• For  $g \cdot z$  large enough all particles die, i.e.  $\overline{\mu}_t \longrightarrow \delta_{\emptyset}$  as  $t \to \infty$ . Equivalently  $k_{\overline{\mu}_t}^{(n)} \longrightarrow 0$  for  $n \ge 1$  as  $t \to \infty$ . Free branching with birth-and-death rates

$$d(x, \gamma \setminus x, \omega) = m + \sum_{y \in \gamma \setminus x} a^{-}(x - y) + g_0 \sum_{y \in \omega} d(x - y)$$
$$b(x, \gamma, \omega) = \sum_{y \in \gamma} a^{+}(x - y) + g_1 \sum_{y \in \omega} b(x - y)$$

- $a^{-}(x y) = a^{-}(y x) \ge 0$  integrable and bounded, competition kernel for cells.
- b(x − y) = b(y − x) ≥ 0 integrable and bounded, proliferation kernel from environment.
- $g_0, g_1 \ge 0$  coupling constant for interaction with environment.

#### Spatial logistic model

### Reduced description

Suppose the following conditions:

- There exists  $\Theta > 0$  such that  $\Theta a^- a^+$  is a stable potential.
- There exists c > 0 such that  $b \le c \cdot d$ .
- *m* is sufficiently large.

Scaling Markov operator  $L_{\gamma}^{S}(\omega) + \frac{1}{\varepsilon}L_{\omega}^{E}$  for  $\varepsilon > 0$  yields when  $\varepsilon \to 0$  reduced description:

$$\overline{d}(x,\gamma\backslash x) = m + \overline{g}(x) + \sum_{y \in \gamma\backslash x} a^{-}(x-y)$$
$$\overline{b}(x,\gamma) = \sum_{y \in \gamma} a^{+}(x-y) + \overline{z}(x)$$

where  $\overline{z}(x) = g_1 \int\limits_{\Gamma^E} \sum\limits_{y \in \omega} b(x - y) \mathrm{d} \mu_{\mathrm{inv}}(\omega)$  and

$$\overline{g}(x) = g_0 \int\limits_{\Gamma^E} \sum_{y \in \omega} d(x - y) \mathrm{d}\mu_{\mathrm{inv}}(\omega)$$

Without presence of environment:

• Dynamics is asymptotically degenerated, i.e.  $\mu_t \longrightarrow \delta_{\emptyset}$  as  $t \to \infty$ . Equivalently  $k_{\mu t}^{(n)} \longrightarrow 0$  for all  $n \ge 1$  as  $t \to \infty$ .

In the presence of environment

• Dynamics has non-trivial invariant measure  $\mu_{\infty}$  such that  $\mu_t \longrightarrow \mu_{\infty}$  as  $t \to \infty$ . Equivalently  $k_{\mu_t}^{(n,m)} \longrightarrow k_{\mu_{\infty}}^{(n,m)}$  for all  $n, m \ge 0$  as  $t \to \infty$ .

After reduced description

• Dynamics has non-trivial invariant measure  $\overline{\mu}_{\infty}$  such that  $\overline{\mu}_t \longrightarrow \overline{\mu}_{\infty}$  as  $t \to \infty$ . Equivalently  $k_{\overline{\mu}_t}^{(n)} \longrightarrow k_{\overline{\mu}_{\infty}}^{(n)}$  for all  $n \ge 0$  as  $t \to \infty$ .

## Consequences

Without presence of environment:

• Kinetic equation is

$$\frac{\partial \rho_t(x)}{\partial t} = -m\rho_t(x) - \int\limits_{\mathbb{R}^d} a^-(x-y)\rho_t(y) \mathrm{d}y\rho_t(x) + \int\limits_{\mathbb{R}^d} a^+(x-y)\rho_t(y) \mathrm{d}y.$$

In the presence of environment

$$\begin{split} \frac{\partial \rho_t^E(x)}{\partial t} &= -\rho_t^E(x) + z e^{-\int\limits_{\mathbb{R}^d} \varphi(x-y)\rho_t^E(y)\mathrm{d}y} \\ \frac{\partial \rho_t^S(x)}{\partial t} &= -m\rho_t^S(x) - \int\limits_{\mathbb{R}^d} a^-(x-y)\rho_t^S(y)\mathrm{d}y\rho_t^S(x) - \int\limits_{\mathbb{R}^d} d(x-y)\rho_t^E(y)\mathrm{d}y\rho_t^S(x) \\ &+ \int\limits_{\mathbb{R}^d} a^+(x-y)\rho_t^S(y)\mathrm{d}y + \int\limits_{\mathbb{R}^d} b(x-y)\rho_t^E(y)\mathrm{d}y. \end{split}$$

After reduced description

$$\frac{\partial \overline{\rho}_t(x)}{\partial t} = -(m + \overline{g}(x))\overline{\rho}_t(x) + \int\limits_{\mathbb{R}^d} a^+(x - y)\overline{\rho}_t(y) \mathrm{d}y + \overline{z}(x).$$

# Dynamics

### Tumour cells

$$d(x, \gamma \setminus x, \omega) = \sum_{y \in \omega} a^{-}(x - y)$$
  
 $b(x, \gamma, \omega) = \sum_{y \in \gamma} a^{+}(x - y)$ 

Immune system

$$d^{E}(x,\gamma,\omega\backslash x) = m + \sum_{y\in\gamma} b^{-}(x-y)$$
$$b^{E}(x,\gamma,\omega) = \sum_{y\in\omega} \left(1 - e^{-E_{\varphi}(y,\gamma)}\right) b^{+}(x-y) + z$$

Derive kinetic equations.

# Dynamics

Kinetic equations

$$\begin{aligned} \frac{\partial \rho_t^{\mathsf{E}}(x)}{\partial t} &= -\int\limits_{\mathbb{R}^d} a^- (x-y) \rho_t^{\mathsf{E}}(y) \mathrm{d}y \rho_t^{\mathsf{S}}(x) + \int\limits_{\mathbb{R}^d} a^+ (x-y) \rho_t^{\mathsf{S}}(y) \mathrm{d}y \\ \frac{\partial \rho_t^{\mathsf{E}}(x)}{\partial t} &= -\left(m - \int\limits_{\mathbb{R}^d} b^- (x-y) \rho_t^{\mathsf{S}}(y) \mathrm{d}y\right) \rho_t^{\mathsf{E}}(x) \\ &+ \int\limits_{\mathbb{R}^d} b^+ (x-y) \left(1 - e^{-\int\limits_{\mathbb{R}^d} \varphi(w-y) \rho_t^{\mathsf{S}}(w) \mathrm{d}w}\right) \rho_t^{\mathsf{E}}(y) \mathrm{d}y + z \end{aligned}$$

Space-homogeneous version:  $X = \rho^{S}$  and  $Y = \rho^{E}$ 

$$X' = (a^+ - a^- Y)X$$
  
 $Y' = z - mY + b^+ Y(1 - e^{-\varphi X}) - b^- XY.$ 

## **Dynamics**

Tumour cells

$$egin{aligned} d(x,\gammaackslash x,\omega) &= m^{\mathcal{S}} + \sum_{y\in\gammaackslash x} b^{-}(x-y) + \sum_{y\in\omega} arphi^{-}(x-y) \ b(x,\gamma,\omega) &= \sum_{y\in\gamma} b^{+}(x-y) + \sum_{y\in\omega} arphi^{+}(x-y) \end{aligned}$$

Immune system

$$egin{aligned} d^{E}(x,\gamma,\omegaackslash x)&=m^{E}+\sum_{y\in\omegaackslash x}a^{-}(x-y)\ b^{E}(x,\gamma,\omega)&=\sum_{y\in\omega}a^{+}(x-y)+z \end{aligned}$$

Environment has still invariant measure  $\mu_{inv}$  with  $\mu_t^{\mathcal{E}} \longrightarrow \mu_{inv}$  as before.

# Reduced description

### Have birth-and-death rates

$$\overline{d}(x,\gamma\backslash x) = m^{S} + \overline{\varphi}^{-}(x) + \sum_{y \in \gamma\backslash x} b^{-}(x-y)$$
$$\overline{b}(x,\gamma) = \sum_{y \in \gamma} b^{+}(x-y) + \overline{z}(x)$$

where 
$$\overline{\varphi}^{-}(x) = \int_{\Gamma^{E}} \sum_{y \in \omega} \varphi^{-}(x - y) d\mu_{inv}(\omega)$$
 and

$$\overline{z}(x) = \int_{\Gamma^E} \sum_{y \in \omega} \varphi^+(x - y) \mathrm{d}\mu_{\mathrm{inv}}(\omega).$$

# Kinetic equations

Without environment

$$\frac{\partial \rho_t(x)}{\partial t} = -\left(m^{\mathsf{s}} + \int\limits_{\mathbb{R}^d} b^-(x-y)\rho_t(y)\mathrm{d}y\right)\rho_t(x) + \int\limits_{\mathbb{R}^d} b^+(x-y)\rho_t(y)\mathrm{d}y.$$

Reduced description

$$\begin{split} \frac{\partial \overline{\rho}_t(x)}{\partial t} &= -\left(m^{\mathsf{S}} + \overline{\varphi}^-(x) + \int\limits_{\mathbb{R}^d} b^-(x-y)\overline{\rho}_t(y) \mathrm{d}y\right) \overline{\rho}_t(x) \\ &+ \int\limits_{\mathbb{R}^d} b^+(x-y)\overline{\rho}_t(y) \mathrm{d}y + \overline{z}(x) \end{split}$$

# Kinetic equations

Theorem

Thank You!